

"INSTABILITY OF SOLUTE PARTICLES"

By

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Abstract:

The thermosolutal stability is investigated by taking the fluid to be viscous, electrically conducting and heterogeneous ,confined between two free boundaries under the action of vertical magnetic field. The study of thermosolutal convection in porous medium in a heterogeneous fluid is of great importance. It finds many applications in biomechanics . It is found that in this case principle of exchange of stabilities is valid and thus the marginal state is a stationary state. Both the suspended particles and the permeability of the medium have destabilizing character. It is found that in this case the marginal state is oscillatory in character and the stable solute gradient has stabilizing effect. The effect of suspended particles and thermal Rayleigh number greater than or equal to solute Rayleigh number is also studied. Further, the effect of suspended particles is to destabilize the layer.



Keywords: Rayleigh number, solute, biomechanics, oscillatory character.

We consider the stability of an incompressible, electrically conducting, viscous, density stratified fluid layer in a porous medium in the presence of a uniform magnetic field and let fluid is rotating about axes $(0, 0, \wedge)$ with solute concentration X. The fluid is to be taken statically non-homogeneous confined between two horizontal boundaries and heated from below. Let T_0 and T_1 ($T_0 \square T_1$) denote the uniform temperatures at the lower and upper boundaries, respectively. we consider the marginal state of the system. We thus take $\square_r \square 0$ and so express $\square \square i\square_1, \square_1$ is real and represent the oscillatory character of the perturbations. we

have

$$\frac{Ra^{2}}{A_{0} \cdot iP_{r}D_{1}} \Box A_{0}(A_{1} \cdot iD_{1}) \cdot \frac{Qa_{x}^{2}A_{0}}{A_{0} \cdot iP_{1}D_{1}} \cdot i\frac{R_{2}a^{2}}{P_{r}D_{1}} \cdot \frac{R_{1}a^{2}}{DA_{0} \cdot iP_{r}D_{1}} \cdot \frac{TD^{2}}{[(A_{0} \cdot iP_{1}D_{1})(A_{1} \cdot iD_{1}) \cdot Qa_{x}^{2}]}$$

Taking the real and imaginary parts of the above equation, we have the real part as:

$$\frac{Ra^{2}A_{0}}{(A_{0}^{2} \bullet P_{r}^{2}\mathcal{D}_{1}^{2})} \Box A_{0}A_{1} \bullet \frac{Qa_{x}^{2}A_{0}}{(A_{0}^{2} \bullet P_{1}^{2}\mathcal{D}_{1}^{2})} \bullet \frac{R_{1}a^{2}\mathcal{D}A_{0}}{(\mathcal{D}^{2}A_{0}^{2} \bullet P_{r}^{2}\mathcal{D}_{1}^{2})} \\ \cdot \frac{T\mathcal{D}^{2}(A_{0}A_{1} \bullet Qa_{x}^{2} \bullet P_{1}\mathcal{D}_{1}^{2})}{(A_{0}A_{1} \bullet Qa_{x}^{2} \bullet P_{1}\mathcal{D}_{1}^{2})^{2} \bullet \mathcal{D}_{1}^{2}(A_{0} \bullet A_{1}P_{1})^{2}}$$



and the imaginary part:

After eliminating the Rayleigh number R from above equations, we find the frequency of oscillation at the marginal state and the Rayleigh number is given and its minimum value is called the critical Rayleigh number corresponding with critical wave number a_C and the frequency \overline{D}_C .

Let us say the non-oscillatory modes exist for which \mathcal{D}_1 is zero and $\mathcal{D} \square \mathcal{D}_2, \mathcal{D}_2$ is real. Hence substituting $\mathcal{D} \square \mathcal{D}_2$ in equation, we get

 $B_0 \mathcal{I}_2^7 \bullet B_1 \mathcal{I}_2^6 \bullet B_2 \mathcal{I}_2^5 \bullet B_3 \mathcal{I}_2^4 \bullet B_4 \mathcal{I}_2^3 \bullet B_5 \mathcal{I}_2^2 \bullet B_6 \mathcal{I}_2 \bullet B_7 \square 0$

where ,

$$B_{0} \square A_{0} P_{r}^{3} P_{1}^{2}$$

$$B_{1} \square A_{0} P_{r}^{2} P_{1} [A_{0} P_{1}(1 \cdot \square \cdot 2P_{r}) \cdot 2P_{r}(A_{0} \cdot P_{1}B)]$$

$$B_{2} \square A_{0} P_{r}^{3} \square P_{1}^{2}(\square A_{0} \cdot P_{r}^{2} A_{1}^{2}) \cdot 2P_{1} P_{r} \square A_{0}(1 \cdot \square) \frac{(A_{0} \cdot P_{1}A_{1})}{3}$$

$$\cdot P_{r}(2A_{0}A_{1} \cdot Qa_{x}^{2}) \square \cdot A_{0}^{2} P_{r}^{2} \square \cdot P_{1}^{2} P_{r}^{2}(R_{1} \cdot R_{2} \cdot R)a^{2} \cdot T\square^{2} P_{1} P_{r}^{3}$$



$$\begin{split} & \mathcal{B}_{3} \, \Big[\qquad \mathcal{A}_{0} P_{r} \{ \mathcal{A}_{0} P_{r}(1 \cdot \overline{D}) [\mathcal{A}_{0}^{2} \cdot P_{1}^{2} \mathcal{A}_{r}^{2} \cdot 2P_{1}(2\mathcal{A}_{0}\mathcal{A} \cdot Qa_{x}^{2})] \\ & \cdot 2\overline{D\mathcal{A}_{0}^{2}} P(\mathcal{A}_{0} \cdot P_{1}\mathcal{A}) \cdot 2P_{r}^{2} (\mathcal{A}_{0}\mathcal{A} + Qa_{x}^{2}) (\mathcal{A}_{0} \cdot P_{1}\mathcal{A}) \} \\ & \cdot R_{2} a^{2} P_{1} P_{1} [2\mathcal{A}_{0} P_{r} \cdot P_{1}(\mathcal{A}_{0}(1 \cdot \overline{D}) \cdot \mathcal{A} P_{r})] \\ & \cdot R_{1} a^{2} P_{1} [P_{1} P_{r}(\mathcal{A}_{0} \cdot \mathcal{A}) \cdot \mathcal{A}_{0}(P_{r} \cdot \overline{D} P_{1}^{2})] \cdot T\overline{D}^{2} P_{r}^{2} \mathcal{A}_{0}(P_{r} \\ & \cdot (1 \cdot \overline{D}) P_{1}) \cdot Ra^{2} P_{1} [P_{1} P_{r}(\mathcal{A}_{0} \cdot \mathcal{A}) \cdot \mathcal{A}_{0}(P_{r} \cdot \overline{D} P_{1}^{2})] \\ B_{4} \, \left[\begin{array}{c} \mathcal{A}_{0} P_{r} \{ P_{r}^{2}(\mathcal{A}_{0}\mathcal{A} - Qa_{x}^{2})^{2} \cdot \overline{D} \mathcal{A}_{0}^{2} [\mathcal{A}_{0}^{2} + P_{1}^{2}\mathcal{A}^{2} \cdot 2P_{1}(Qa_{x}^{2} \cdot 2\mathcal{A}_{0}\mathcal{A})] \\ & \cdot (1 \cdot \overline{D}) P_{1} \cdot Ra^{2} P_{1} [P_{1} P_{r}(\mathcal{A}_{0} \cdot \mathcal{A}) + \mathcal{A}_{0}(P_{r} \cdot \overline{D} P_{1}^{2})] \\ B_{4} \, \left[\begin{array}{c} \mathcal{A}_{0} \mathcal{A}_{r} (P_{x}^{2}(\mathcal{A}_{0}\mathcal{A} - Qa_{x}^{2})^{2} \cdot \overline{D} \mathcal{A}_{0}^{2} [\mathcal{A}_{0}^{2} + P_{1}^{2}\mathcal{A}^{2} \cdot 2P_{1}(Qa_{x}^{2} \cdot 2\mathcal{A}_{0}\mathcal{A})] \\ & \cdot (1 \cdot \overline{D}) P_{1} \cdot (P_{r}^{2}(\mathcal{A}_{0}\mathcal{A} - Qa_{x}^{2})^{2} \cdot \overline{D} \mathcal{A}_{0}^{2} [\mathcal{A}_{0}^{2} + P_{1}^{2}\mathcal{A}^{2} \cdot 2P_{1}(Qa_{x}^{2} \cdot 2\mathcal{A}_{0}\mathcal{A})] \\ & \cdot (\mathcal{A}_{0} \mathcal{A}_{0} Qa_{x}^{2})(\mathcal{A}_{0} + P_{1}\mathcal{A}) \mathcal{A}_{0} P_{1}(1 \cdot \overline{D}) \} R_{2} a^{2} \{\mathcal{A}_{0}^{2}(\mathcal{D} P_{r} + P_{r} \cdot DP_{1}^{2}) \\ & \cdot (\mathcal{A}_{0}\mathcal{A} + Qa_{x}^{2})(\mathcal{A}_{0} + P_{1}\mathcal{A}) \mathcal{A}_{0} P_{1}^{2} P_{1}\mathcal{A}_{0}^{2} \mathcal{A}_{0}^{2} \mathcal{A}_$$

 $A_1 \square A_0 \bullet B$

(i) Since $R_2 \square 0$ then $B_7 \square 0$, and the product of all seven roots is positive, therefore at least one root is positive, these making the system is unstable. Therefore, we conclude that non-oscillatory modes are unstable.



(ii) If $R_2 \square 0$, then $B_7 \square 0$ and product of all roots is negative, hence either all roots are negative or at least one root is positive, thus both negative or positive roots are possible.

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